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Inverse Approaches to Drying of Thin Bodies With Significant Shrinkage Effects

This paper deals with the application of inverse concepts to the drying of bodies that undergo changes in their dimensions. Simultaneous estimation is performed of moisture diffusivity, together with the thermal conductivity, heat capacity, density, and phase conversion factor of a drying body, as well as the heat and mass transfer coefficients and the relative humidity of drying air. This was accomplished by using only temperature measurements. A mathematical model of the drying process of shrinking bodies has been developed where the moisture content and temperature fields in the drying body are expressed by a system of two coupled partial differential equations. The shrinkage effect was incorporated through the experimentally obtained changes of the specific volume of the drying body in an experimental convective dryer. The proposed method was applied to the process of drying potatoes. For the estimation of the unknown parameters, the transient readings of a single temperature sensor located in the midplane of the potato slice, exposed to convective drying, have been used. The Levenberg–Marquardt method and a hybrid optimization method of minimization of the least-squares norm are used to solve the present parameter estimation problem. Analyses of the sensitivity coefficients and of the determinant of the information matrix are presented as well.

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Keywords: inverse approach, drying, thermophysical properties, heat and mass transfer coefficients

1 Introduction

2 There are several methods for describing the direct problem of
3 complex simultaneous heat and moisture transport processes
4 within a drying material. In the approach proposed by Luikov [1]
5 the moisture and temperature fields in the drying body are ex-
6 pressed by a system of two coupled partial differential equations.
7 The system of equations incorporates coefficients that must be
8 determined experimentally.
9 All the coefficients, except for the moisture diffusivity, can be
10 relatively easily determined by experiments [2,3]. A number of
11 methods for the experimental determination of the moisture diffu-
12 sivity exist such as: sorption kinetics methods, permeation meth-
13 ods, concentration-distance methods, drying methods, radiotracer
14 methods, and methods based on the techniques of electron spin
15 resonance and nuclear magnetic resonance, but there is no stan-
16 dard method. The adoption of a generalized method for moisture
17 diffusivity estimation would be of great importance.
18 We have recently analyzed a method for moisture diffusivity
19 estimation by the temperature response of a drying body [4–11].

The main idea of this method is to make use of the interrelation
between the heat and mass transport processes within the drying
body and from its surface to the surroundings. Then, the moisture
diffusivity can be estimated on the basis of an accurate and easy to
perform single thermocouple temperature measurement by using
an inverse approach.

The objective of this paper is an analysis of the possibility of
the simultaneous estimation of the moisture diffusivity, together
with other thermophysical properties of vegetables, as well as the
heat and mass transfer coefficients. The method requires a single
drying experiment and a single temperature measurement probe.
As a representative drying vegetable product, thin slices of potato
have been chosen. An analysis of the influence of the drying air
velocity, temperature and relative humidity, drying body dimen-
sions, and drying time on the moisture diffusivity estimation, en-
ables the design of appropriate experiments to be conducted as
well. In order to realize this analysis, the sensitivity coefficients
and the determinant of the information matrix were calculated for
the characteristic drying regimes and drying body dimensions.

Physical Problem and Mathematical Formulation

The physical problem involves a single slice of a potato of
thickness $2L$ initially at uniform temperature and uniform mois-
ture content (Fig. 1). The surfaces of the drying body are in con-
tact with the drying air, thus resulting in a convective boundary
condition for both the temperature and the moisture content.

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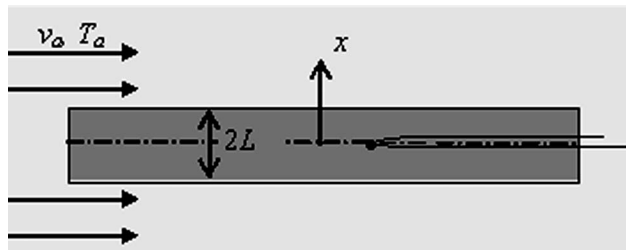


Fig. 1 Scheme of the drying experiment

45 problem is symmetrical relative to the mid-plane of the slice. The
 46 thickness of the body changes during the drying from $2L_0$ to $2L_f$.
 47 In the case of an infinite flat plate the unsteady temperature,
 48 $T(x, t)$, and moisture content, $X(x, t)$, fields in the drying body are
 49 expressed by the following system of coupled nonlinear partial
 50 differential equations for energy and moisture transport

$$51 \quad c\rho_s \frac{\partial T}{\partial t} = \frac{\partial}{\partial x} \left(k \frac{\partial T}{\partial x} \right) + \varepsilon \Delta H \frac{\partial(\rho_s X)}{\partial t} \quad (1)$$

$$52 \quad \frac{\partial(\rho_s X)}{\partial t} = \frac{\partial}{\partial x} \left(D\rho_s \frac{\partial X}{\partial x} + D\rho_s \delta \frac{\partial T}{\partial x} \right) \quad (2)$$

53 Here, t , x , c , k , ΔH , ε , δ , D , and ρ_s are time, normal distance
 54 from the midplane of the plate, specific heat, thermal conductivity,
 55 latent heat of vaporization, ratio of water evaporation rate to the
 56 reduction rate of the moisture content, thermogradient coefficient,
 57 moisture diffusivity, and density of dry solid, respectively.

58 From the experimental and numerical examinations of the trans-
 59 sient moisture and temperature profiles [12] it was concluded that
 60 for practical calculations, the influence of any thermodiffusion is
 61 small and can be ignored. Consequently, $\delta=0$ was utilized in this
 62 paper.

63 The shrinkage effect of the drying body was incorporated
 64 through the changes of the specific volume of the drying body.
 65 There are several models for describing the changes of the specific
 66 volume of the body during drying. In this paper, linear relation-
 67 ship between the specific volume, ν_s , and the moisture content, X ,
 68 has been used

$$69 \quad \nu_s = \frac{1}{\rho_s} = \frac{V}{m_s} = \frac{1 + \beta' X}{\rho_{b0}} \quad (3)$$

70 Here, m_s is the mass of the dry material of the drying body, V is
 71 the volume of the drying body, ρ_{b0} is the density of a fully dried
 72 body, and β' is the shrinkage coefficient.

73 Substituting the above expression for $\rho_s (=1/\nu_s)$ into Eqs. (1)
 74 and (2) and rearranging with $\delta=0$, results in

$$75 \quad \frac{\partial T}{\partial t} = \frac{k}{\rho_s c} \frac{\partial^2 T}{\partial x^2} + \frac{\varepsilon \Delta H}{c} \frac{\rho_s}{\rho_{b0}} \frac{\partial X}{\partial t} \quad (4)$$

$$76 \quad \frac{\partial X}{\partial t} = D \frac{\rho_{b0}}{\rho_s} \frac{\partial^2 X}{\partial x^2} + \frac{\rho_{b0}}{\rho_s^2} \frac{\partial(D\rho_s)}{\partial x} \frac{\partial X}{\partial x} \quad (5)$$

77 The problem of the moving boundaries due to the changes of
 78 the dimensions of the body during the drying was resolved by
 79 introducing the dimensionless coordinate

$$80 \quad \psi = \frac{x}{L(t)} \quad (6)$$

81 Consequently, the resulting system of equations for the tem-
 82 perature and moisture content prediction becomes

$$83 \quad \frac{\partial T}{\partial t} = \frac{k}{\rho_s c L^2} \frac{\partial^2 T}{\partial \psi^2} + \frac{\psi dL}{L dt} \frac{\partial T}{\partial \psi} + \frac{\varepsilon \Delta H}{c} \frac{\rho_s}{\rho_{b0}} \left(\frac{\partial X}{\partial t} - \frac{\psi dL}{L dt} \frac{\partial X}{\partial \psi} \right) \quad (7)$$

$$\frac{\partial X}{\partial t} = D \frac{\rho_{b0}}{\rho_s} \frac{1}{L^2} \frac{\partial^2 X}{\partial \psi^2} + \left[\frac{\rho_{b0}}{\rho_s^2} \frac{1}{L^2} \frac{\partial(D\rho_s)}{\partial \psi} + \frac{\psi dL}{L dt} \right] \frac{\partial X}{\partial \psi} \quad (8) \quad 84$$

The initial conditions are 85

$$t = 0: \quad T(\psi, 0) = T_0, \quad X(\psi, 0) = X_0 \quad (9) \quad 86$$

The temperature and the moisture content boundary conditions 87
 on the surfaces of the drying body in contact with the drying air 88
 are 89

$$-k \frac{1}{L} \left(\frac{\partial T}{\partial \psi} \right)_{\psi=1} + j_q - \Delta H(1 - \varepsilon) j_m = 0 \quad 90$$

$$D\rho_s \frac{1}{L} \left(\frac{\partial X}{\partial \psi} \right)_{\psi=1} + j_m = 0 \quad (10) \quad 91$$

The convective heat flux, $j_q(t)$, and mass flux, $j_m(t)$, on these 92
 surfaces are 93

$$j_q = h(T_a - T_{x=L}) \quad 94$$

$$j_m = h_D(C_{x=L} - C_a) \quad (11) \quad 95$$

where h is the heat transfer coefficient, and h_D is the mass transfer 96
 coefficient, T_a is the temperature of the drying air, and $T_{x=L}$ is the 97
 temperature on the surfaces of the drying body. The water vapor 98
 concentration in the drying air, C_a , is calculated from 99

$$C_a = \frac{\varphi p_s(T_a)}{R_w T_{k,a}} \quad (12) \quad 100$$

where φ is the relative humidity of the drying air and p_s is the 101
 saturation pressure. The water vapor concentration of the air in 102
 equilibrium with the surface of the body exposed to convection is 103
 calculated from 104

$$C_{x=L} = \frac{a(T_{x=L}, X_{x=L}) p_s(T_{x=L})}{R_w T_{k,x=L}} \quad (13) \quad 105$$

The water activity, a , or the equilibrium relative humidity of the 106
 air in contact with the convection surface at temperature $T_{x=L}$ and 107
 moisture content $X_{x=L}$ are calculated from experimental water 108
 sorption isotherms. 109

The boundary conditions on the midplane of the drying slice are 110

$$\left(\frac{\partial T}{\partial \psi} \right)_{\psi=0} = 0, \quad \left(\frac{\partial X}{\partial \psi} \right)_{\psi=0} = 0 \quad (14) \quad 111$$

Problem defined by Eqs. (7)–(14) is referred to as a direct prob- 112
 lem when initial and boundary conditions as well as all the param- 113
 eters appearing in the formulation are known. The objective of 114
 the direct problem is to determine the temperature and moisture 115
 content fields in the drying body. 116

In order to approximate the solution of Eqs. (7) and (8), an 117
 explicit numerical procedure has been used. 118

The Drying Body Properties 119

In this paper, application of the proposed method for the esti- 120
 mation of the thermophysical properties of vegetables has been 121
 analyzed. As a representative vegetable product, a potato was chosen. 122
123

Heat capacity of food materials can be taken as equal to the 124
 sum of the heat capacity of solid matter and water absorbed by 125
 that solid 126

$$c = c_s + c_w X \quad (15) \quad 127$$

Although the heat capacity of solid matter, c_s , and water, c_w , are 128
 functions of the temperature, constant values have been most 129
 widely used. 130

131 From Ref. [12] it was also concluded that for practical calcula-
 132 tions the system of the two simultaneous partial differential equa-
 133 tions could be used by treating the thermal conductivity, k , and the
 134 phase conversion factor, ε , as constants.
 135 Moisture diffusivity of foods is a function of the temperature
 136 and the moisture content as well. The moisture diffusivity depen-
 137 dence of the moisture content for a potato is not clearly expressed
 138 [[13], p. 216], and it is very often considered as an Arrhenius-type
 139 temperature function [14,15]

$$140 \quad D = D_0 \exp[-E_0/(RT_k)] \quad (16)$$

141 with constant values of the Arrhenius factor, D_0 , and the activa-
 142 tion energy for moisture diffusion, E_0 .

143 The variation in water activity with change in moisture content
 144 of samples at a specified temperature is defined by sorption iso-
 145 therms. There are many different models for describing the sorp-
 146 tion isotherms of foods [3]. In recent years, the most widely ac-
 147 cepted and efficient model for sorption isotherms of foods has
 148 been the Guggenheim–Anderson–de Boer (GAB) model

$$149 \quad X = \frac{X_m CKa}{(1 - Ka)(1 - Ka + CKa)} \quad (17)$$

150 The monolayer moisture, X_m , and the adsorption constants C and
 151 K are related as Arrhenius type equations with the Arrhenius fac-
 152 tors X_{m0} , C_0 , and K_0 , and the energy terms ΔH_X , ΔH_C , and ΔH_K ,
 153 respectively.

154 Inverse Approach

155 The inverse problem in this paper is solved as a parameters
 156 estimation approach. The estimation methodology used is based
 157 on the minimization of the ordinary least square norm

$$158 \quad E(\mathbf{P}) = [\mathbf{Y} - \mathbf{T}(\mathbf{P})]^T [\mathbf{Y} - \mathbf{T}(\mathbf{P})] \quad (18)$$

159 Here, $\mathbf{Y}^T = [Y_1, Y_2, \dots, Y_{\text{imax}}]$ is the vector of measured tem-
 160 peratures, $\mathbf{T}^T = [T_1(\mathbf{P}), T_2(\mathbf{P}), \dots, T_{\text{imax}}(\mathbf{P})]$ is the vector of esti-
 161 mated temperatures at time t_i ($i=1, 2, \dots, \text{imax}$), \mathbf{P}^T
 162 $= [P_1, P_2, \dots, P_N]$ is the vector of unknown parameters, imax is the
 163 total number of measurements, and N is the total number of un-
 164 known parameters ($\text{imax} \geq N$).

165 A hybrid optimization algorithm OPTRAN [16] and the
 166 Levenberg–Marquardt method [17–19] have been utilized for the
 167 minimization of $E(\mathbf{P})$ representing the solution of the present pa-
 168 rameter estimation problem.

169 The Levenberg–Marquardt method is a quite stable, powerful,
 170 and straightforward gradient search minimization algorithm that
 171 has been applied to a variety of inverse problems. It belongs to a
 172 general class of damped least square methods. The solution for
 173 vector \mathbf{P} is achieved using the following iterative procedure

$$174 \quad \mathbf{P}^{r+1} = \mathbf{P}^r + [(\mathbf{J}^r)^T \mathbf{J}^r + \mu \mathbf{I}]^{-1} (\mathbf{J}^r)^T [\mathbf{Y} - \mathbf{T}(\mathbf{P}^r)] \quad (19)$$

175 where r is the number of iterations, \mathbf{I} is identity matrix, μ is the
 176 damping parameter, and \mathbf{J} is the sensitivity matrix defined as

$$177 \quad \mathbf{J} = \begin{bmatrix} \frac{\partial T_1}{\partial P_1} & \dots & \frac{\partial T_1}{\partial P_N} \\ \vdots & & \\ \frac{\partial T_{\text{imax}}}{\partial P_1} & \dots & \frac{\partial T_{\text{imax}}}{\partial P_N} \end{bmatrix} \quad (20)$$

178 Near the initial guess, the problem is generally ill conditioned
 179 so that large damping parameters are chosen thus making the term
 180 $\mu \mathbf{I}$ large as compared to term $\mathbf{J}^T \mathbf{J}$. The term $\mu \mathbf{I}$ damps instabilities
 181 due to the ill-conditioned character of the problem. So, the matrix
 182 $\mathbf{J}^T \mathbf{J}$ is not required to be nonsingular at the beginning of the itera-
 183 tions and the procedure tends towards a slow-convergent steepest
 184 descent method. As the iteration process approaches the con-
 185 verged solution, the damping parameter decreases, and the
 186 Levenberg–Marquardt method tends towards a Gauss method. In

fact, this method is a compromise between the steepest descent 187
 and Gauss method by choosing μ so as to follow the Gauss 188
 method to as large an extent as possible, while retaining a bias 189
 towards the steepest descent direction to prevent instabilities. The 190
 presented iterative procedure terminates if the norm of gradient of 191
 $E(\mathbf{P})$ is sufficiently small, if the ratio of the norm of the gradient 192
 of $E(\mathbf{P})$ to $E(\mathbf{P})$ is small enough, or if the changes in the vector of 193
 parameters are very small. 194

195 An alternative to the Levenberg–Marquardt algorithm, espe-
 cially when searching for a global optimum of a function with
 possible multiple minima, is the hybrid optimization program OP-
 TRAN [16]. OPTRAN incorporates six of the most popular optimi-
 zation algorithms: the Davidon–Fletcher–Powell gradient search 199
 [20], sequential quadratic programming algorithm [21], 200
 Pshenichny–Danilin quasi-Newtonian algorithm [22], a modified 201
 Nelder–Mead simplex algorithm [23], a genetic algorithm [24], 202
 and a differential evolution algorithm [25]. Each algorithm pro- 203
 vides a unique approach to optimization with varying degrees of 204
 convergence, reliability, and robustness at different stages during 205
 the iterative optimization procedure. The hybrid optimizer OPTRAN 206
 includes a set of rules and switching criteria to automatically 207
 switch back and forth among the different algorithms as the itera- 208
 tive process proceeds in order to avoid local minima and accelera- 209
 te convergence towards a global minimum. 210

211 The population matrix was updated every iteration with new
 designs and ranked according to the value of the objective func-
 tion, in this case the ordinary least square norm. As the optimiza-
 tion process proceeded, the population evolved towards its global
 minimum. The optimization problem was completed when one of
 several stopping criteria was achieved: the maximum number of
 iterations or objective function evaluations was exceeded, the best
 design in the population was equivalent to a target design, or the
 optimization program tried all six algorithms, but failed to pro-
 duce a nonnegligible decrease in the objective function. The last
 criterion usually indicated that a global minimum had been found. 221

222 Parameters Estimation Analysis

223 For the inverse problem of interest here, the moisture diffusiv-
 ity parameters, together with other thermophysical properties of
 the potato as well as the heat and mass transfer coefficients and
 the relative humidity of the drying air, are treated as unknown
 226 parameters. 227

228 Thus, in the inverse problem the analyzed vector of the un-
 known parameters was 229

$$230 \quad \mathbf{P}^T = [D_0, E_0, \rho_s, c_s, k, \varepsilon, h, h_D, \varphi] \quad (21)$$

231 For the simultaneous estimation of these unknown parameters,
 232 the transient reading of a single temperature sensor located at the
 position $x=0$, has been considered. 233

234 The possibility of the simultaneous estimation of the
 temperature-dependent moisture diffusivity together with the
 other thermophysical properties of the potato as well as the heat
 and mass transfer coefficients and the relative humidity of the
 drying air depends on the boundary conditions and dimensions of
 the drying sample. An analysis of the influence of the drying air
 parameters and dimensions of the drying sample needed for the
 design of the appropriate experiment have been conducted in this
 paper. In order to perform this analysis, the sensitivity coefficients
 242 have been calculated. 243

244 The sensitivity coefficients analysis has been carried out for an
 infinite flat plate model of a slice of potato with initial moisture
 content of $X(x, 0) = 5.00$ kg/kg and initial temperature $T(x, 0)$
 246 $= 20.0^\circ\text{C}$. The drying air bulk temperature, T_a , was varied be-
 247 tween 40 and 80°C , the convection heat transfer coefficient be-
 248 tween 27 and 33 $\text{W m}^{-2} \text{K}^{-1}$, and the initial thickness, $2L_0$, of the
 249 potato slice between 2 and 6 mm. 250

251 From the sensitivity coefficients analysis the following experi-
 mental parameters were chosen: $T_a = 60^\circ\text{C}$, $2L_0 = 3$ mm, h
 252 $= 30$ $\text{W m}^{-2} \text{K}^{-1}$, $h_D = 3.36 \times 10^{-2}$ m s^{-1} , and $\varphi = 0.09$. Figure 2 253

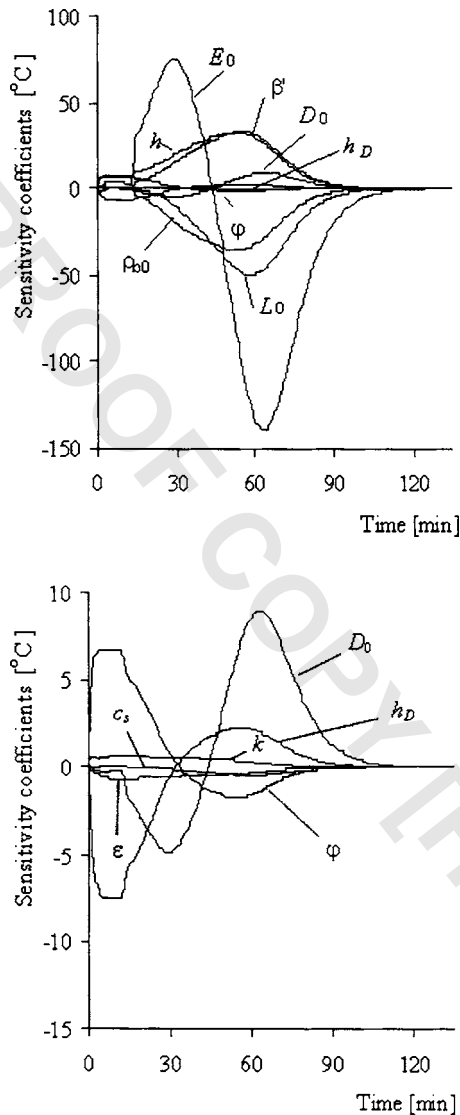


Fig. 2 Relative sensitivity coefficients for temperature

254 shows the relative sensitivity coefficients $P_j \partial T_i / \partial P_j$, i
 255 = 1, 2, ..., imax, for temperature with respect to the unknown pa-
 256 rameters, for this case.
 257 It can be seen that the relative sensitivity coefficients with re-
 258 spect to the phase conversion factor, ϵ , and the thermal conduc-
 259 tivity, k , are very small. This indicates that ϵ and k cannot be
 260 estimated in this case. This also indicates that the influence of the
 261 phase conversion factor and the thermal conductivity on the tran-
 262 sient moisture content and temperature profiles is very small in
 263 this case. This can be explained by the very small heat transfer
 264 Biot number ($Bi = hL/k \leq hL_0/k = 0.11$) and consequently very
 265 small temperature gradients inside the body during the drying pro-
 266 cess. For these reasons, the phase conversion factor and the ther-
 267 mal conductivity were treated as known quantities for the exami-
 268 nation described below.
 269 The heat capacity of wet potato was taken as equal to the sum
 270 of the heat capacity of solid matter and absorbed water, Eq. (15).
 271 Since the heat capacity of the solid matter, c_s , presents only a few
 272 percent of the overall heat capacity of the potato, the relative
 273 sensitivity coefficients with respect to the heat capacity of solid
 274 matter is also very small. Consequently, the value of the heat
 275 capacity of the solid matter was also taken as known.
 276 The relative sensitivity coefficients with respect to the density
 277 of the fully dried body, ρ_{b0} , and the shrinkage coefficient, β' , are

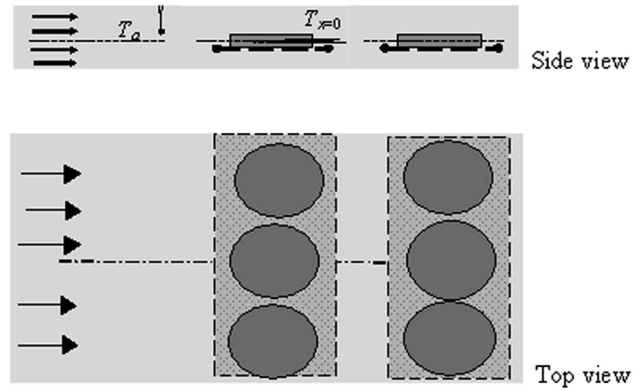


Fig. 3 Scheme of the experimental arrangement

relatively high. Despite this, because the shrinkage effect of the
 drying body was incorporated through the changes of the specific
 volume of the drying body, these parameters were determined by
 separate experiments.

The relative sensitivity coefficients with respect to the initial
 potato slice thickness are high as well, but the initial slice thick-
 ness was measured with sufficient accuracy, so it is also taken as
 a known parameter.

It can be seen that the temperature sensitivity coefficient with
 respect to the convection mass transfer coefficient h_D is very small
 relative to the temperature sensitivity coefficient with respect to
 the convection heat transfer coefficient, h . The very high mass
 transfer Biot number and the very small heat transfer Biot number
 can explain this. To overcome this problem, in this paper the mass
 transfer coefficient was related to the heat transfer coefficient
 through the analogy between the heat and mass transfer processes
 in the boundary layer over a drying body [10]

$$h_D = 0.95 \frac{D_a}{k_a} h \quad (22)$$

where D_a and k_a are the moisture diffusivity and thermal conduc-
 tivity in the air, respectively. The obtained relation is very close to
 the well-known Lewis relation. By using the above relation be-
 tween the heat and mass transfer coefficients, they can be esti-
 mated simultaneously, so that only the heat transfer coefficient is
 regarded as an unknown parameter.

Experiment

Real experiments have been conducted to investigate the appli-
 cability of the method to food processing, when involving drying
 of thin flat samples. The experiments have been conducted on the
 experimental setup that is designed to imitate an industrial con-
 vective dryer.

Drying of approximately three millimetres thick potato slices
 have been examined. The slices have been in contact with the
 drying air from the top and the bottom surfaces. Two shelves,
 each holding three moist potato slices have been introduced into
 the rectangular experimental channel of dimensions 25 × 200 mm.
 A microthermocouple was inserted in the midplane, ($x=0$), of
 each of the three slices on the first shelf. An arithmetical mean
 of the readings from the three thermocouples was used as a
 transient temperature reading, ($T_{x=0}$), for the estimation of the
 unknown parameters. The potato slices on the second shelf were
 weighed every ten minutes in order to obtain the volume-averaged
 moisture content change during drying. The temperature of the
 drying air, T_a , has been recorded as well. The initial moisture
 content, X_0 , and the initial potato slices thickness, $2L_0$, were
 measured for each of the experiments.

The change of the specific volume of the drying body was
 determined by a separate experiment. The cylindrical potato slices

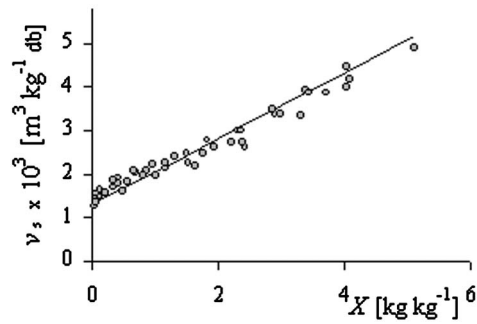


Fig. 4 Change of the specific volume during the drying of potato slices

325 with diameter of approximately 40 mm and thickness of approxi-
 326 mately 3 mm have been placed on the second shelf and dried until
 327 the equilibrium moisture content has been reached. The dimen-
 328 sions and the mass of the slices were measured every 10 min. The
 329 initial moisture content and the initial potato slices thicknesses
 330 were measured as well. The experiment was repeated for different
 331 temperatures and speed of the drying air. The drying air tempera-
 332 ture was varied between 50 and 70°C, and the drying air speed
 333 between 1.0 and 3.0 m s⁻¹.

334 The relative errors of the measurements were estimated be-
 335 tween 0.1% and 1.0% for the mass and 0.3–2.5% for the dimen-
 336 sions of the slices. Microthermocouples were calibrated, relative
 337 to each other, within 0.2°C in the range of 20–80°C.

338 **Results and Discussion**

339 From the parameters estimation analysis it was concluded that
 340 the moisture diffusivity parameters, D_0 and E_0 , the convection
 341 heat and mass transfer coefficients, h and h_D , and the relative
 342 humidity of the drying air, φ , will be treated as unknown param-
 343 eters in this paper. All other quantities appearing in the direct
 344 problem formulation were taken as known.

345 Our experimental results for the changes of the specific volume
 346 of drying potato slices, (Fig. 4), confirm the expression (3) with
 347 $\rho_{b0}=755 \text{ kg m}^{-3}$ and the shrinkage coefficient $\beta^t=0.57$. The heat
 348 capacity was calculated from Eq. (15). The following values, pro-
 349 posed in Ref. [26] for potatoes, were used: $c_s=1381 \text{ J kg}^{-1} \text{ K}^{-1}$
 350 and $c_w=4187 \text{ J kg}^{-1} \text{ K}^{-1}$. A mean value of $k=0.40 \text{ W m}^{-1} \text{ K}^{-1}$
 351 from the results obtained in Ref. [27] for the thermal conductivity
 352 of potato was utilized in this paper. The influence of the phase
 353 conversion factor ($0 \leq \varepsilon \leq 1$) on the transient moisture content and
 354 temperature profiles is very small. A mean value of $\varepsilon=0.5$ was
 355 used in the paper. For the GAB isotherm equation parameters the
 356 Gane experimental results for potatoes [[3], p. 45] were used in
 357 this paper ($C_0=6.609 \times 10^{-1}$; $\Delta H_C=528.4 \text{ kJ kg}^{-1}$; $K_0=0.606$;
 358 $\Delta H_K=53.33 \text{ kJ kg}^{-1}$; $\Delta H_X=123.6 \text{ kJ kg}^{-1}$), except for X_{m0} . The
 359 value of $X_{m0}=3.8 \times 10^{-2}$ was obtained from our experimental re-
 360 sults.

361 A number of drying experiments with similar experimental
 362 conditions, ($T_a=56.6\text{--}59.5^\circ\text{C}$, $2L_0=2.36\text{--}3.14 \text{ mm}$, $X_0=3.70\text{--}$
 363 4.83 kg/kg , and $T_0=14.9\text{--}17.7^\circ\text{C}$), have been carried out.

364 The experimental drying time was estimated from the determi-
 365 nant of the information matrix. Figure 5 presents the transient

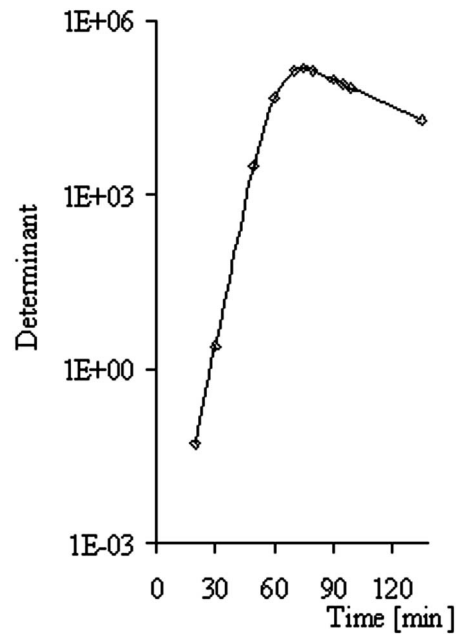


Fig. 5 Determinant of the information matrix

variation of the determinant of the information matrix if D_0 , E_0 , h , h_D , and φ are simultaneously considered as unknown. Elements of this determinant of the information matrix were defined [19] for a large, but fixed number of transient temperature measurements (451 in this case).

The maximum determinant value corresponds to the drying time when near equilibrium moisture content and temperature profiles have been reached.

The relative temperature sensitivity coefficients with respect to the moisture diffusivity parameters, D_0 and E_0 , are almost linearly dependent, (Fig. 2). Despite this, we were able to obtain results using the OPTRAN [16] and the Levenberg–Marquardt algorithm [18]. Table 1 shows the computationally obtained parameters and rms error for experiment E1: $T_a=58.13^\circ\text{C}$, $2L_0=3.14 \text{ mm}$, $X_0=4.80 \text{ kg/kg}$, and $T_0=17.53^\circ\text{C}$. The rms changes and the convergence of the estimated values of the unknown parameters to the final values during the iterative process of the Levenberg–Marquardt method, for experiment E1 are shown in Fig. 6.

In Fig. 7 the estimated moisture diffusivities are compared with the results published by other authors that used different methods.

In Fig. 8 the experimental transient temperature reading, $T_{x=0}$, and the experimental volume-averaged moisture content change during drying are compared with numerical solutions for the estimated parameters. Very good agreements were obtained. The temperature changes during the weighing of the slices on the second shelf (every 10 min the second shelf together with the slices was taken outside the channel for 15 s to be weighed) can be seen in Fig. 8.

In Fig. 9, the experimental transient temperature reading, $T_{x=0}$, and the experimental volume-averaged moisture content changes during drying are compared with the numerical solutions with the estimated parameters in the cases when the shrinkage effect was

Table 1 Estimated parameters and rms error

	$D_0 \cdot 10^3$ (m s ⁻¹)	E_0 (kJ mol ⁻¹)	h (W m ⁻² K ⁻¹)	$h_D \cdot 10^2$ (m s ⁻¹)	φ (–)	rms (°C)
Initial guess	0.25	35	25	2.8	0.125	4.93
Estimated values	7.985	43.3	31.08	3.48	0.0899	0.55

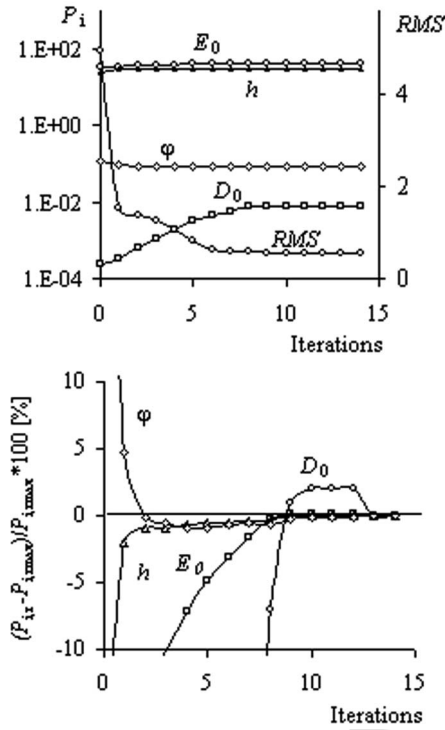


Fig. 6 Convergence history of rms errors and estimated parameters

398 incorporated and when it was not incorporated. It is very clear that
 399 the shrinkage effect cannot be ignored in the calculations of the
 400 drying processes of potato slices.

401 **Conclusions**

402 The inverse problem of simultaneous estimation of thermo-
 403 physical properties and the boundary condition parameters of drying
 404 thin slices of vegetables by using only temperature measurements
 405 has been analysed. For this, a mathematical model of drying of
 406 shrinking bodies has been developed. As a representative vegetable
 407 product, a slice of a potato has been chosen.

408 It can be concluded that in the convective drying experiment it
 409 is possible, based on a single thermocouple temperature response,

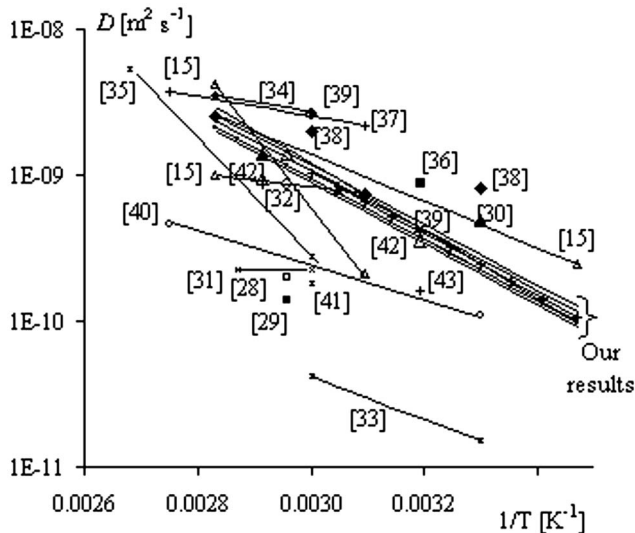
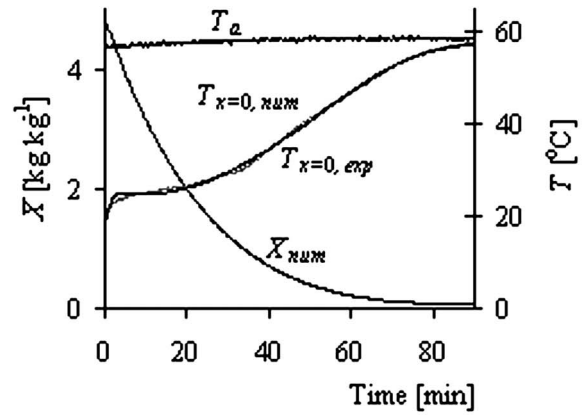
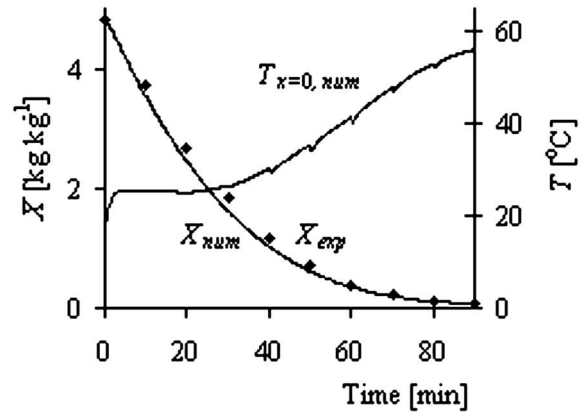


Fig. 7 Moisture diffusivity of potatoes



(a)



(b)

Fig. 8 Time-variations during drying: The midplane temperature, $T_{x=0}$, the temperature of the drying air, T_a , and the volume-averaged moisture content, X : (a) the first shelf, and (b) the second shelf

to estimate simultaneously the two moisture diffusivity parameters, the convection heat and mass transfer coefficients, and the relative humidity of the drying air.

Estimated moisture diffusivities compare well with the values obtained by other authors who utilized different methods.

Very good agreement between the experimental and numerical temperature and volume-averaged moisture content changes during drying has been obtained.

Since the relative temperature sensitivity coefficients with respect to the moisture diffusivity parameters in the Arrhenius-type

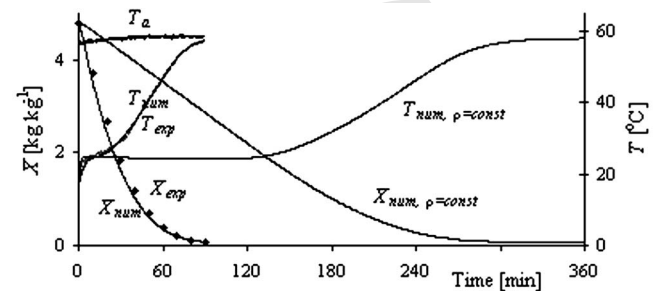


Fig. 9 Changes during drying with shrinkage effect and without shrinkage effect: The midplane temperature, $T_{x=0}$ and the volume-averaged moisture content, X

420 function are linearly dependent, other models for describing the
421 moisture content and temperature-dependent moisture diffusivity
422 could be analyzed in the future [28–43].

423 Nomenclature

424 a = water activity
425 c = heat capacity (dry basis), $\text{J K}^{-1} \text{kg}^{-1} \text{db}$
426 C = concentration of water vapor, kg m^{-3}
427 D = moisture diffusivity, $\text{m}^2 \text{s}^{-1}$
428 D_0 = Arrhenius factor, $\text{m}^2 \text{s}^{-1}$
429 E_0 = activation energy J kg^{-1}
430 E = ordinary least square norm $(^\circ\text{C})^2$
431 h = heat transfer coefficient, $\text{W m}^{-2} \text{K}^{-1}$
432 h_D = mass transfer coefficient, m s^{-1}
433 ΔH = latent heat of vaporization, J kg^{-1}
434 \mathbf{I} = identity matrix
435 j_m = mass flux, $\text{kg m}^{-2} \text{s}^{-1}$
436 j_q = heat flux, W m^{-2}
437 \mathbf{J} = sensitivity matrix
438 k = thermal conductivity, $\text{W m}^{-1} \text{K}^{-1}$
439 L = flat plate thickness, m
440 m = mass, kg
441 p_s = saturation pressure, Pa
442 \mathbf{P} = vector of unknown parameter
443 R = absolute gas constant, $\text{J K}^{-1} \text{mol}^{-1}$
444 R_w = specific gas constant, $\text{J K}^{-1} \text{kg}^{-1}$
445 t = time, s
446 T = temperature, $^\circ\text{C}$
447 T_k = temperature, K
448 \mathbf{T} = vector of estimated temperature, $^\circ\text{C}$
449 v = specific volume, $\text{m}^3 \text{kg}^{-1}$
450 V = volume, m^3
451 x = distance from the midplane, m
452 X = moisture content (dry basis), $\text{kg kg}^{-1} \text{db}$
453 \mathbf{Y} = vector of measured temperature, $^\circ\text{C}$

474 Greek symbols

475 β' = shrinkage coefficient, -
476 δ = thermo-gradient coefficient, K^{-1}
477 ε = phase conversion factor
478 μ = damping parameter
479 ρ = density, kg m^{-3}
480 φ = relative humidity
481 ψ = dimensionless coordinate

484 Subscripts

485 a = drying air
486 b_0 = fully dried body
487 f = final
488 m = monolayer
489 w = water
490 s = dry solid

491 Superscripts

492 r = number of iterations
493 T = transposed

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